

207(2): Metric Compatibility and Equation of Motion

The metric compatibility equation is:

$$D_\rho g_{\mu\nu} = 0 \quad - (1)$$

where $g_{00} = 1, g_{11} = - \left(1 + \frac{r^2}{f^2} \right) - (2)$

and $f = \frac{dr}{d\theta} - (3)$

Let: $\rho = 0 - (4)$

then: $D_0 g_{\mu\nu} - \Gamma^\lambda_{\mu 0} g_{\lambda\nu} - \Gamma^\lambda_{\nu 0} g_{\mu\lambda} = 0 - (5)$

so $D_0 g_{00} - \Gamma^\lambda_{00} g_{\lambda 0} - \Gamma^\lambda_{00} g_{0\lambda} = 0 - (6)$

However: $\Gamma^\lambda_{00} = 0, - (7)$

so $\boxed{D_0 g_{00} = 0,} - (8)$

which is true because of eq. (2).

Also: $D_0 g_{11} - \Gamma^\lambda_{01} g_{\lambda 1} - \Gamma^\lambda_{01} g_{1\lambda} = 0 - (9)$

i.e $D_0 g_{11} - \Gamma^1_{01} g_{11} - \Gamma^1_{01} g_{11} = 0 - (10)$

and $\boxed{\Gamma^1_{01} = \frac{1}{2g_{11}} D_0 g_{11}} - (11)$

Let $\rho = 1 - (12)$

2) then:

$$\partial_1 g_{\mu\nu} - \Gamma_{1\mu}^{\lambda} g_{\lambda\nu} - \Gamma_{1\nu}^{\lambda} g_{\mu\lambda} = 0 \quad - (13)$$

then $\partial_1 g_{00} - \Gamma_{10}^{\lambda} g_{\lambda 0} - \Gamma_{10}^{\lambda} g_{0\lambda} = 0 \quad - (14)$

$$2\Gamma_{10}^0 g_{00} = 0 \quad - (15)$$

$$\boxed{\Gamma_{10}^0 = 0} \quad - (16)$$

Also: $\partial_1 g_{11} - \Gamma_{11}^{\lambda} g_{\lambda 1} - \Gamma_{11}^{\lambda} g_{1\lambda} = 0 \quad - (17)$

i.e. $\boxed{\partial_1 g_{11} = 0} \quad - (18)$

Let $p=2$, then

$$\partial_2 g_{\mu\nu} - \Gamma_{2\mu}^{\lambda} g_{\lambda\nu} - \Gamma_{2\nu}^{\lambda} g_{\mu\lambda} = 0 \quad - (19)$$

For example:

$$\partial_2 g_{00} - \Gamma_{20}^{\lambda} g_{\lambda 0} - \Gamma_{20}^{\lambda} g_{0\lambda} = 0 \quad - (20)$$

i.e. $2\Gamma_{20}^0 g_{00} = 0 \quad - (21)$

$$\boxed{\Gamma_{20}^0 = 0} \quad - (22)$$

Also $\partial_2 g_{11} - \Gamma_{21}^{\lambda} g_{\lambda 1} - \Gamma_{21}^{\lambda} g_{1\lambda} = 0 \quad - (23)$

and $\boxed{\Gamma_{21}^1 = \frac{1}{2g_{11}} \partial_2 g_{11}} \quad - (24)$

3) The time equation of motion for Q. denoted by using output 3b are:

$$6(1+f) \frac{d^2 f}{dt^2} = 5 \left(\frac{df}{dt} \right)^2 \quad - (25)$$

and

$$\frac{d}{dt} \left(3 \frac{df}{dr} - (1+f) \frac{df}{dr} \right) = 0$$

$$\therefore \frac{d}{dt} \left((3 - (1+f)) \frac{df}{dr} \right) = 0 \quad - (26)$$

For the precessing ellipse:

$$r = \frac{d}{1 + e \cos(x\theta)} \quad - (27)$$

$$\left. \begin{aligned} f &= \frac{dr}{d\theta} = \left(\frac{x e}{d} \right) r^2 \sin(x\theta) \\ &= \left(\frac{x e}{d} \right) r^2 \left(1 - \frac{1}{e^2} \left(\frac{d}{r} - 1 \right)^2 \right)^{1/2} \end{aligned} \right\} \quad - (28)$$

Denote $g(r(t)) = 3 - (1+f) \frac{df}{dr} \quad - (29)$

then eq. (26) is:

$$\boxed{\frac{dg(r(t))}{dt} = 0} \quad - (30)$$

4) From eq. (2a):

$$g = g(r, \theta, t) \quad - (31)$$

so:
$$\frac{dg}{dt} = \frac{\partial g}{\partial t} + \frac{\partial g}{\partial r} \frac{dr}{dt} + \frac{\partial g}{\partial \theta} \frac{d\theta}{dt} \quad - (32)$$

so eq. (30) is:
$$\frac{dg}{dt} = \frac{\partial g}{\partial r} \frac{dr}{dt} + \frac{\partial g}{\partial \theta} \frac{d\theta}{dt} \quad - (33)$$

because:
$$\frac{\partial g}{\partial r} = 0 \quad - (34)$$

However, if:
$$g = g(r, \theta) \quad - (35)$$

eq. (33) is always true.

Therefore the identity (26) is always true
if $f = f(r, \theta)$, or more accurately:

$$f = f(r(t), \theta(t)) \quad - (36)$$

Now use:
$$\frac{df}{dt} = \frac{df}{dr} \frac{dr}{dt} \quad - (37)$$

and
$$\frac{dr}{dt} = \frac{dr}{d\theta} \frac{d\theta}{dt} \quad - (38)$$

From note 207(i):

$$\frac{dr}{d\theta} = 0 \quad - (39)$$

5) so the identity (25) is always true because both sides go to infinity.

This analysis shows yet again that the
Einstein identity is always true. It is

$$D_{\mu} T^{\mu\nu} = R^{\mu\nu}{}_{\mu}{}^{\nu} \quad - (40)$$

and is an example of the contracted identity.