

230(5): Development of the General Fermi equation.

The general fermi equation can be written as:

$$p^a p_\mu = m^2 c^2 g_\mu^a \quad - (1)$$

where

$$p_\mu = p_0 g_\mu^a \quad - (2)$$

is the momentum tetrad. So:

$$p^a p_\mu = \left(\frac{m^2 c^2}{p_0} \right) p_\mu^a \quad - (3)$$

The LCC minimal prescription is:

$$p_\mu^a \rightarrow p_\mu^a - e A_\mu^a \quad - (4)$$

so in the presence of electromagnetic radiation:

$$p^a p_\mu = \left(\frac{m^2 c^2}{p_0} \right) (p_\mu^a - e A_\mu^a) \quad - (5)$$

$$\text{i.e. } p_\mu^a - e A_\mu^a = \left(\frac{p_0}{m^2 c^2} \right) p^a p_\mu \quad - (6)$$

$$= \left(\frac{p_0}{m^2 c^2} \right) g_\mu^a p^a p_\mu \quad - (7)$$

$$\text{Therefore: } g_\mu^a (p_\mu^a - e A_\mu^a) = \left(\frac{p_0}{m^2 c^2} \right) p^a p_\mu \quad - (8)$$

2) In this equation:

$$\gamma_a^4 p_\mu^a = p_0 ; e \gamma_a^4 A_\mu^a = e A_0 - (9)$$

$$\text{So } \boxed{p_\mu^4 = m^2 c^2 \left(1 - \frac{e A_0}{p_0} \right)} \quad - (10)$$

In this equation A_0 may be regarded as the magnitude of a vacuum potential.

Eq. (10) may be written as:

$$E^2 - c^2 p^2 = m^2 c^4 \left(1 - \frac{e A_0}{p_0} \right) \quad - (11)$$

$$E^2 - c^2 p^2 = m^2 c^4 \left(1 - \frac{e A_0}{p_0} \right) \quad - (12)$$

i.e.

$$E - mc^2 \left(1 - \frac{e A_0}{p_0} \right)^{1/2} = \frac{c^2 p^2}{E + mc^2 \left(1 - \frac{e A_0}{p_0} \right)^{1/2}}$$

This equation may now be approximated and reduced to a Schrodinger type equation.

In the non relativistic limit on the right hand

side:

$$E \rightarrow mc^2 \quad - (13)$$

For a weak perturbation:

$$e A_0 \ll p_0 \quad - (14)$$

3) then:

$$E - mc^2 \left(1 - \frac{eA_0}{p_0}\right)^{1/2} \sim \frac{p^2}{2m}, \quad - (15)$$

$$\text{and } \left(1 - \frac{eA_0}{p_0}\right)^{1/2} \sim 1 - \frac{eA_0}{2p_0}, \quad - (16)$$

$$\text{so } E - mc^2 + \frac{emc^2 A_0}{2p_0} \sim \frac{p^2}{2m} \quad - (17)$$

$$\text{Denote } E_1 = E - mc^2 \quad - (18)$$

$$\text{then } E_1 + \frac{emc^2 A_0}{2p_0} \sim \frac{p^2}{2m} \quad - (19)$$

- (20)

This quantizes to:

$$\boxed{-\frac{\hbar^2 \nabla^2}{2m} \psi = \left(E_1 + \frac{emc^2 A_0}{2p_0}\right) \psi}$$

In presence of a potential V :

$$\left(-\frac{\hbar^2 \nabla^2}{2m} + V\right) \psi = \left(E_1 + \frac{emc^2 A_0}{2p_0}\right) \psi \quad - (21)$$

$$\text{i.e. } -\frac{\hbar^2 \nabla^2}{2m} \psi = \left(E_1 - V + \frac{emc^2 A_0}{2p_0}\right) \psi \quad - (22)$$

4) So:

$$\nabla^2 \phi = -\frac{2m}{\hbar^2} \left(E_1 - V + \frac{emc^2 A_0}{2p_0} \right) \phi \quad - (23)$$

This equation may now be used in tunnelling
theory.
