

322(4): The Gravitomagnetic Field in General Dynamics

Consider the position coordinate:

$$\underline{r} = x \underline{i} + y \underline{j} + z \underline{k} \quad - (1)$$

In cylindrical polar coordinates:

$$x = r \cos \theta, \quad y = r \sin \theta, \quad - (2)$$

$$\underline{r} = r \cos \theta \underline{i} + r \sin \theta \underline{j} + z \underline{k} \quad - (3)$$

so:

$$= \dot{r} \underline{e}_r + \dot{z} \underline{k}$$

With reference to "Vector Analysis Problem Solver",
problem 21-12, p. 1029, the velocity is:

$$\underline{v} = \dot{r} \underline{e}_r + \omega r \underline{e}_\theta + \dot{z} \underline{k} \quad - (4)$$

and the acceleration is:

$$\underline{a} = (\ddot{r} - r\dot{\theta}^2) \underline{e}_r + (r\ddot{\theta} + 2\dot{r}\dot{\theta}) \underline{e}_\theta + \ddot{z} \underline{k} \quad - (5)$$

where by definition:

$$\underline{\omega} = \frac{d\theta}{dt} \underline{k} \quad - (6)$$

Therefore the general gravitomagnetic field is:

$$\underline{\Omega} = -\frac{1}{2c} \underline{v} \times \underline{a} \quad - (7)$$

$$= -\frac{1}{2c} \begin{vmatrix} \underline{e}_r & \underline{e}_\theta & \underline{k} \\ \dot{r} & \omega r & \dot{z} \\ (\ddot{r} - r\dot{\theta}^2) & (r\ddot{\theta} + 2\dot{r}\dot{\theta}) & \ddot{z} \end{vmatrix}$$

2) In general, this expression contains the three Newtonian forces contained in eq. (5).

Three Dimensional Orbits

These are defined by the central force:

$$\underline{F} = -\frac{mMG}{r^2} \underline{e}_r \quad - (8)$$

The Lagrangian is:

$$L = \frac{1}{2} m v^2 - U \quad - (9)$$

where:

$$v^2 = \dot{r}^2 + \dot{\theta}^2 r^2 + \dot{z}^2 \quad - (10)$$

and

$$U = -\frac{MG}{r} \quad - (11)$$

The three Euler Lagrange equations are:

$$\frac{\partial L}{\partial \theta} = \frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} \quad - (12)$$

$$\frac{\partial L}{\partial r} = \frac{d}{dt} \frac{\partial L}{\partial \dot{r}} \quad - (13)$$

$$\frac{\partial L}{\partial z} = \frac{d}{dt} \frac{\partial L}{\partial \dot{z}} \quad - (14)$$

Eq. (12) gives:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} = 0 \quad - (15)$$

3) Eq. (13) gives:

$$F(r) = -\frac{\partial U}{\partial r} = m(\ddot{r} - r\dot{\theta}^2) \quad - (16)$$

Eq. (14) gives:

$$\frac{dZ}{dt} = 0 \quad - (17)$$

Eq. (15) is conservation of angular momentum, eq. (16) is the 1689 Leibniz equation of orbits, and eq. (17) shows that Z does not change with time.

The angular momentum is defined by:

$$\underline{L} = m \underline{r} \times \underline{v} \quad - (18)$$

$$= m \begin{vmatrix} \underline{e}_r & \underline{e}_\theta & \underline{k} \\ r & 0 & Z \\ \dot{r} & \omega r & 0 \end{vmatrix}$$

Therefore:

$$\underline{L} = m(\dot{r}Z \underline{e}_\theta + \omega r^2 \underline{k}) \quad - (19)$$

an equation which shows clearly that \underline{L} is not perpendicular to the orbital plane.

From eq. (19):

4)

$$L_z = m r^2 \omega - (20)$$

$$L_\theta = m r \dot{\theta} - (21)$$

It follows that:

$$L_z = \frac{\partial \mathcal{L}}{\partial \dot{\theta}} = m r^2 \omega - (22)$$

and

$$\boxed{\omega_z = \frac{L_z}{m r^2} = \frac{d\theta}{dt}} - (23)$$

The Lagrangian analysis shows that:

$$\frac{dL_z}{dt} = 0 - (24)$$

Therefore L_z is a constant of motion.

Note carefully that:

$$\frac{dL}{dt} \neq 0 - (25)$$

where

$$L^2 = L_\theta^2 + L_z^2 - (26)$$

From eqs. (16) and (23) it follows as in Maria and Thoma pages 249 ff of the 2nd ed. that:

$$5) \quad F(r) = -\frac{L_z^2}{mr^3} \left(\frac{d^2}{d\theta^2} \left(\frac{1}{r} \right) + \frac{1}{r} \right) - (27)$$

which is the Binet equation of orbit

By definition:

$$\omega_z = \dot{\theta} = \frac{L_z}{mr^2} \quad - (28)$$

so

$$\ddot{\theta} = \frac{d\dot{\theta}}{dr} \frac{dr}{dt} = \dot{r} \frac{d}{dr} \left(\frac{L_z}{mr^2} \right)$$

$$= -\frac{2\dot{r}L_z}{mr^3} \quad - (29)$$

Therefore:

$$r\ddot{\theta} = -\frac{2\dot{r}L_z}{mr^2} \quad - (30)$$

Similarly:

$$2\dot{r}\dot{\theta} = \frac{2\dot{r}L_z}{mr^2} \quad - (31)$$

It follows that:

$$r\ddot{\theta} + 2\dot{r}\dot{\theta} = 0 \quad - (32)$$

From eqs. (4), (5), (17) and (32):

$$\underline{v} = \dot{r} \underline{e}_r + \omega r \underline{e}_\theta \quad - (33)$$

$$\underline{a} = (\ddot{r} - r\dot{\theta}^2) \underline{e}_r \quad - (34)$$

b) and: $\underline{r} = r \underline{e}_r + Z \underline{k} \quad - (35)$
 $\underline{L} = m (i Z \underline{e}_\theta + \omega r^2 \underline{k}) \quad - (36)$

The difference between planar and free dimensional orbital theory is that in planar

theory: $\underline{r} = r \underline{e}_r \quad - (37)$
 and $\underline{L} = m \omega r^2 \underline{k} \quad - (38)$

In planar theory: $Z = 0 \quad - (39)$

For clarity eq. (35) can be defined as:

$$\underline{r}_{\text{total}} = r \underline{e}_r + Z \underline{k} \quad - (40)$$

so $r_{\text{total}}^2 = r^2 + Z^2 \quad - (41)$

In general, the observed orbit is:

$$r = \frac{d}{1 + \epsilon \cos(x\theta)} \quad - (42)$$

where $r = (r_{\text{total}}^2 - Z^2)^{1/2} \quad - (43)$

and where $L_z^2 = m^2 M G d \quad - (44)$

The angular velocity, gravitomagnetic field,

) and the mass current remains as in Note 322(2):

$$\underline{\omega} = \frac{L_z}{mr^2} \underline{k} \quad - (45)$$

$$\underline{a} = - \left(\frac{MG}{rc^2} \right) \frac{1}{r^3} \underline{k} \quad - (46)$$

$$\underline{J}_n = - \frac{3M}{4\pi n} \frac{L_z}{r^4} \underline{e}_\theta \quad - (47)$$

with L_z taking the place of L .

Finally, note that:

$$\begin{aligned} \underline{\omega} \times \underline{r} &= \omega_z \underline{k} \times (r \underline{e}_r + z \underline{k}) \\ &= \omega_z r \underline{k} \times \underline{e}_r \quad - (48) \\ &= \omega_z r \underline{e}_\theta \end{aligned}$$

Therefore:

$$\underline{F} = m \underline{g} - \underline{\omega} \times (\underline{\omega} \times \underline{r}) = - \frac{mMG \underline{e}_r}{r^2} \quad - (49)$$

The added information is a 3-D orbit is:

$$r_{\text{total}}^2 = \left(\frac{d}{1 + \epsilon \cos(\chi\theta)} \right)^2 + z^2 \quad - (50)$$

9) In the most general dynamics:

$$J = \frac{1}{2} m (\dot{r}^2 + \dot{\theta}^2 r^2 + \dot{z}^2) - U(r, \theta, z) \quad - (51)$$

and in the most general dynamics:

$$\underline{a} = \ddot{r} \underline{e}_r - \underline{\omega} \times (\underline{\omega} \times \underline{r}) + \frac{d\underline{\omega}}{dt} \times \underline{r} + 2 \underline{\omega} \times \frac{d\underline{r}}{dt} \underline{e}_r \quad - (52)$$

in which the centripetal acceleration is $-\underline{\omega} \times (\underline{\omega} \times \underline{r})$,
 the Coriolis acceleration is $2 \underline{\omega} \times \frac{d\underline{r}}{dt} \underline{e}_r$ and the
 kind non-Newtonian acceleration is $\frac{d\underline{\omega}}{dt} \times \underline{r}$. The
 velocity in general dynamics is:

$$\underline{v} = \dot{r} \underline{e}_r + \underline{\omega} \times \underline{r} \quad - (53)$$

and the most general gravitomagnetic field is

$$\underline{\Omega} = -\frac{1}{c^2} \underline{v} \times \underline{a} \quad - (54)$$

Therefore $\underline{\Omega}$ must be worked out for eqs (52)
 and (53).
