

47(7): Analysis of the body in Spherical Polar Coordinates

In spherical polar coordinates:

$$\underline{r} = r \underline{e}_r \quad - (1)$$

and the velocity is:

$$\underline{v} = \dot{r} \underline{e}_r + r \dot{\theta} \underline{e}_\theta + r \sin \phi \dot{\phi} \underline{e}_\phi \quad - (2)$$

"Vega Analysis Problem Solver" 21-25.

The acceleration is:

$$\underline{a} = a_r \underline{e}_r + a_\theta \underline{e}_\theta + a_\phi \underline{e}_\phi \quad - (3)$$

where:

$$a_r = \ddot{r} - r \dot{\theta}^2 - r \sin^2 \theta \dot{\phi}^2 \quad - (4)$$

$$a_\theta = 2 \dot{r} \dot{\theta} + r \ddot{\theta} - r \sin \theta \cos \theta \dot{\phi}^2 \quad - (5)$$

$$a_\phi = 2 \dot{r} \dot{\phi} \sin \theta + 2 r \dot{\theta} \dot{\phi} \cos \theta + r \sin \theta \ddot{\phi} \quad - (6)$$

In a central force such as the gravitational force or a spring:

$$\underline{F} = m \frac{d^2 \underline{r}}{dt^2} = m (a_r \underline{e}_r + a_\theta \underline{e}_\theta + a_\phi \underline{e}_\phi) \quad - (7)$$

Therefore for the spring:

$$-\frac{mg}{r} = \ddot{r} - r \dot{\theta}^2 - r \sin^2 \theta \dot{\phi}^2 \quad - (8)$$

$$2 \dot{r} \dot{\theta} + r \ddot{\theta} - r \sin \theta \cos \theta \dot{\phi}^2 = 0 \quad - (9)$$

$$2 \dot{r} \dot{\phi} \sin \theta + 2 r \dot{\theta} \dot{\phi} \cos \theta + r \sin \theta \ddot{\phi} = 0 \quad - (10)$$

3) So the net force on the gyro is:

$$\underline{F} = m \frac{d^2 \underline{r}}{dt^2} = -\frac{mMg}{r^2} \underline{e}_r$$

$$= m(\ddot{r} - r\dot{\theta}^2 - r\sin^2\theta\dot{\phi}^2)\underline{e}_r \quad - (11)$$

i.e. the gyro generates a centrifugal force in three dimensions:

$$\underline{F}(\text{cent}) = mr(\dot{\theta}^2 + \sin^2\theta\dot{\phi}^2)\underline{e}_r \quad - (12)$$

Note that:

$$\frac{d^2 \underline{r}}{dt^2} = (\ddot{r} - r\dot{\theta}^2 - r\sin^2\theta\dot{\phi}^2)\underline{e}_r \quad - (13)$$

s. the presence of the centrifugal force does not change the equation:

$$\underline{F} = m \frac{d^2 \underline{r}}{dt^2} = -\frac{mMg}{r^2} \underline{e}_r \quad - (14)$$

and the gyro can never lift off the ground because it always obeys eq. (14).

Gyroscope dynamics are described by:

$$\ddot{r} = -\frac{Mg}{r^2} + r(\dot{\theta}^2 + \sin^2\theta\dot{\phi}^2) \quad - (15)$$

which is the three dimensional equivalent of
Leibnitz orbital equation:

$$\ddot{r} = -\frac{MG}{r^2} + r\dot{\phi}^2 \quad (16)$$

In spherical polar coordinates:

$$\frac{d\underline{e}_r}{dt} = \dot{\theta} \underline{e}_\theta + \sin\theta \dot{\phi} \underline{e}_\phi \quad (17)$$

("Vector analysis Problem Solver" 21-24). So:

$$\frac{d\underline{e}_r}{dt} \cdot \frac{d\underline{e}_r}{dt} = \dot{\theta}^2 + \sin^2\theta \dot{\phi}^2 \quad (18)$$

From eqs (15) and (18):

$$\ddot{\underline{r}} = -\frac{MG}{r^2} + r \frac{d\underline{e}_r}{dt} \cdot \frac{d\underline{e}_r}{dt} \quad (19)$$

also

$$\underline{r} = r \underline{e}_r \quad (20)$$

Eq. (19) appears to be a new gyroscope equation.

In order to lift the gyroscope an additional force \underline{F}_1 is needed so that:

$$\underline{F} = \underline{F}_{grav} + \underline{F}_1 > \underline{0} \quad (21)$$

In classical dynamics this lift can

1. Is generated by a torque:

$$\underline{\tau}_r = \underline{r} \times \underline{F}_1 \quad - (22)$$

If:

$$\underline{F}_1 = F_1 \underline{e}_r \quad - (23)$$

experimentally the eq. (19) becomes:

$$\ddot{r} = -\frac{MG}{r^2} + r \frac{d\underline{e}_r}{dt} \cdot \frac{d\underline{e}_r}{dt} + F_1 \quad - (24)$$

otherwise \underline{F}_1 can be generated by a conservative derivative:

$$\underline{F}_1 = m(\underline{v} \cdot \underline{\nabla}) \underline{v} \quad - (25).$$
