

### 370(4): General Dynamics of the Orbiting Asymmetric Top

Consider an asymmetric top whose centre of mass is orbiting an object of mass  $M$ . The mass of the asymmetric top is  $m$  and its principal moments of inertia are  $I_1$ ,  $I_2$  and  $I_3$ . The potential energy of attraction between  $m$  and  $M$  is considered to be:

$$U = U(r, \theta, \phi, \psi) \quad - (1)$$

where  $r$  is the distance between  $m$  and  $M$ , and where  $\theta$ ,  $\phi$  and  $\psi$  are the Euler angles defined in the frame  $(1, 2, 3)$  of the principal moments of inertia.

The Lagrangian is:

$$L = \frac{1}{2} m \dot{\underline{r}} \cdot \dot{\underline{r}} + \frac{1}{2} (I_1 \omega_1^2 + I_2 \omega_2^2 + I_3 \omega_3^2) - U(r, \theta, \phi, \psi) \quad - (2)$$

and the Lagrange variables are  $r$ ,  $\theta$ ,  $\phi$  and  $\psi$ . So the following four Euler Lagrange equations must be solved simultaneously:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\underline{r}}} = \frac{\partial L}{\partial \underline{r}} \quad - (3)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} = \frac{\partial L}{\partial \phi} \quad - (4)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} = \frac{\partial L}{\partial \theta} \quad - (5)$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\psi}} = \frac{\partial L}{\partial \psi} \quad - (6)$$

2) Here:

$$\omega_1 = \dot{\phi} \sin \theta \sin \phi + \dot{\theta} \cos \phi \quad - (7)$$

$$\omega_2 = \dot{\phi} \sin \theta \cos \phi - \dot{\theta} \sin \phi \quad - (8)$$

$$\omega_3 = \dot{\phi} \cos \theta + \dot{\psi} \quad - (9)$$

Eq. (3) gives:

$$m \ddot{\underline{r}} = -\nabla U(\underline{r}, \theta, \phi, \psi) \quad - (10)$$

$\underline{L}, \underline{R}$  is the force equation of the system.

We have:

$$\omega_1^2 = (\dot{\phi} \sin \theta \sin \phi + \dot{\theta} \cos \phi)^2 \quad - (11)$$

$$= \dot{\phi}^2 \sin^2 \theta \sin^2 \phi + 2 \dot{\phi} \dot{\theta} \sin \theta \sin \phi \cos \phi + \dot{\theta}^2 \cos^2 \phi \quad - (12)$$

$$\omega_2^2 = (\dot{\phi} \sin \theta \cos \phi - \dot{\theta} \sin \phi)^2$$

$$= \dot{\phi}^2 \sin^2 \theta \cos^2 \phi - 2 \dot{\theta} \dot{\phi} \sin \theta \sin \phi \cos \phi + \dot{\theta}^2 \sin^2 \phi \quad - (13)$$

$$\omega_3^2 = (\dot{\phi} \cos \theta + \dot{\psi})^2 = \dot{\phi}^2 \cos^2 \theta + 2 \dot{\phi} \dot{\psi} \cos \theta + \dot{\psi}^2$$

So:

$$\mathcal{L} = \frac{1}{2} m \dot{\underline{r}} \cdot \dot{\underline{r}} + \frac{1}{2} (I_1 + I_2) (\dot{\phi}^2 \sin^2 \theta + \dot{\theta}^2)$$

$$+ \dot{\phi} \dot{\theta} \sin \theta \sin \phi \cos \phi (I_1 - I_2) + \frac{1}{2} I_3 (\dot{\phi}^2 \cos^2 \theta + 2 \dot{\phi} \dot{\psi} \cos \theta + \dot{\psi}^2)$$

$$- U(\underline{r}, \theta, \phi, \psi) \quad - (14)$$

It follows that:

$$\frac{\partial \mathcal{L}}{\partial \phi} = - \frac{\partial U}{\partial \phi} \quad - (15)$$

$$\frac{\partial \mathcal{L}}{\partial \dot{\phi}} = (\underline{I}_1 + \underline{I}_2) \dot{\phi} \sin^2 \theta + (\underline{I}_1 - \underline{I}_2) \dot{\theta} \sin \theta \sin \psi \cos \psi + \underline{I}_3 (\dot{\phi} \cos^2 \theta + \dot{\psi} \cos \theta) \quad - (16)$$

Eq. (4) gives:

$$(\underline{I}_1 + \underline{I}_2) \dot{\phi} \sin^2 \theta + (\underline{I}_1 - \underline{I}_2) \dot{\theta} \sin \theta \sin \psi \cos \psi + \underline{I}_3 \cos \theta (\dot{\phi} \cos \theta + \dot{\psi}) = - \frac{\partial \mathcal{U}}{\partial \phi} \quad - (17)$$

$$= (\underline{I}_1 + \underline{I}_2) \frac{d}{dt} (\dot{\phi} \sin^2 \theta) + (\underline{I}_1 - \underline{I}_2) \frac{d}{dt} (\dot{\theta} \sin \theta \sin \psi \cos \psi) + \underline{I}_3 \frac{d}{dt} (\cos \theta (\dot{\phi} \cos \theta + \dot{\psi}))$$

Similarly:

$$\frac{\partial \mathcal{L}}{\partial \dot{\psi}} = - \frac{\partial \mathcal{U}}{\partial \psi} + (\underline{I}_1 - \underline{I}_2) \dot{\phi} \dot{\theta} \sin \theta (\cos^2 \psi - \sin^2 \psi) \quad - (18)$$

$$\text{and } \frac{\partial \mathcal{L}}{\partial \dot{\phi}} = \underline{I}_3 (\dot{\phi} \cos \theta + \dot{\psi}) \quad - (19)$$

so:

$$\underline{I}_3 \frac{d}{dt} (\dot{\phi} \cos \theta + \dot{\psi}) - (\underline{I}_1 - \underline{I}_2) \dot{\phi} \dot{\theta} \sin \theta (\cos^2 \psi - \sin^2 \psi) = - \frac{\partial \mathcal{U}}{\partial \psi} \quad - (20)$$

Finally:

$$\frac{\partial \mathcal{L}}{\partial \theta} = (\underline{I}_1 + \underline{I}_2) \sin \theta \cos \theta + \dot{\phi} \dot{\theta} \cos \theta \sin \psi \cos \psi (\underline{I}_1 - \underline{I}_2) - \underline{I}_3 (\dot{\phi}^2 \sin \theta \cos \theta + \dot{\phi} \dot{\psi} \sin \theta) - \frac{\partial \mathcal{U}}{\partial \theta} \quad - (21)$$

4) and

$$\frac{\partial \mathcal{L}}{\partial \dot{\theta}} = (\mathbb{I}_1 + \mathbb{I}_2) \dot{\theta} + (\mathbb{I}_1 - \mathbb{I}_2) \dot{\phi} \sin \theta \sin \phi \cos \phi - (22)$$

So:

$$\begin{aligned} & (\mathbb{I}_1 + \mathbb{I}_2) \ddot{\theta} + (\mathbb{I}_1 - \mathbb{I}_2) \frac{d}{dt} (\dot{\phi} \sin \theta \sin \phi \cos \phi) \\ & - (\mathbb{I}_1 + \mathbb{I}_2) \sin \theta \cos \theta - (\mathbb{I}_1 - \mathbb{I}_2) \dot{\phi} \dot{\theta} \cos \theta \sin \phi \cos \phi \\ & + \mathbb{I}_3 (\dot{\phi}^2 \sin \theta \cos \theta + \dot{\phi} \ddot{\phi} \sin \theta) = - \frac{\partial U}{\partial \theta} \end{aligned} - (23)$$

Therefore eqs. (15), (20), (17) and (23) must be solved simultaneously.

It is clear that the dynamics depend on  $\partial U / \partial \underline{r}$ ,  $\partial U / \partial \theta$ ,  $\partial U / \partial \phi$  and  $\partial U / \partial \dot{\phi}$ . Therefore models for these functions give orbital characteristics via eq. (10), and different patterns of  $\theta(t)$ ,  $\phi(t)$  and  $\dot{\phi}(t)$ .

Example of Modelling the Potential.

Assume:

$$\frac{\partial U}{\partial \phi} = \frac{\partial U}{\partial \dot{\phi}} = 0 - (24)$$

and

$$U = U(\underline{r} / \theta) - (25)$$

it follows that:

$$\frac{\partial U}{\partial \theta} = \frac{\partial U}{\partial r} \frac{\partial r}{\partial \theta} \quad - (26)$$

$$= \frac{\partial r}{\partial \theta} \nabla U$$

The assumption that  $r$  is a function of  $\theta$  introduces dependence of the rotations and precession on  $r$ .

$$|\nabla U| = \frac{nmG}{r^2} = \frac{\partial U}{\partial r} \quad - (27)$$

for

$$\frac{\partial U}{\partial \theta} = \frac{nmG}{r^2} \frac{\partial r}{\partial \theta} \quad - (28)$$

and this can be used in Eq. (23).

In the plane polar system  $(r, \theta_1)$  Eq.

10) gives

$$\ddot{r} - r\dot{\theta}_1^2 = -\frac{\partial U}{\partial r} \quad - (29)$$

$$r\ddot{\theta}_1 + 2\dot{r}\dot{\theta}_1 = 0 \quad - (30)$$

and

The complete system consists of eqs. (15), (17), (20), (23), and (28) to (30).

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