

572(1): The Basic Equations of Lagrangian Quantum Mechanics

With reference to 4FT 371:

$$-\hbar^2 \nabla^2 \psi = p^2 \psi = \frac{L^2}{a} \left(\frac{2}{r} - \frac{1}{a} \right) \psi \quad (1)$$

where

$$\psi = \psi(r) \psi(\theta, \phi) \quad (2)$$

ψ is the wavefunction, a product of radial and angular wavefunctions. The Lagrangian is:

$$L = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\beta}^2) - U \quad (3)$$

where

$$\dot{\beta}^2 = \dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta \quad (4)$$

The proper Lagrange variables are r and β , so we have two Euler Lagrange equations:

$$\frac{\partial L}{\partial r} = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{r}} \right) \quad (5)$$

and

$$\frac{\partial L}{\partial \beta} = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\beta}} \right) \quad (6)$$

For the H atom: $U(r) = -\frac{e^2}{4\pi\epsilon_0 r} \quad (7)$

in which case: $r = \frac{a}{1 + \epsilon \cos \beta} \quad (8)$

and the classical momentum is:

$$p^2 = \frac{L^2}{d} \left(\frac{2}{r} - \frac{1}{a} \right) - (9)$$

The semi major axis of the ellipse (8) is

$$a = \frac{d}{1-e^2} - (10)$$

Defining

$$U = -\frac{k}{r} - (11)$$

then

$$k = \frac{e^2}{4\pi\epsilon_0} - (12)$$

The half right latitude is:

$$d = \frac{L^2}{mk} = \left(\frac{4\pi\epsilon_0}{e^2} \right) \frac{L^2}{m} - (13)$$

which is Eq. (15) of Note 371(8), Q.E.D.

Eq. (13) can be expressed as:

$$d = \frac{L^2}{mhc\alpha_{fs}} - (14)$$

where the fine structure constant is:

$$\alpha_{fs} = \frac{e^2}{4\pi\hbar c\epsilon_0} - (15)$$

The eccentricity of the three dimensional ellipse (8) is defined as:

$$e^2 = 1 + \frac{2EL^2}{mk^2} - (16)$$

where the total energy of the H atom is:

$$E = -\frac{me^4}{32\pi^2 \epsilon_0^2 \hbar^2 n^2} \quad (17)$$

where n is the principal quantum number.

Therefore:

$$\epsilon^2 = 1 - \frac{1}{n^2} \left(\frac{L}{\hbar} \right)^2 \quad (18)$$

The angular momentum quantizes as:

$$L^2 \psi = l(l+1) \hbar^2 \psi \quad (19)$$

where l is the angular momentum quantum number. So the classical L^2 is the expectation value:

$$\begin{aligned} L^2 &= \langle L^2 \rangle = \int \psi^* L^2 \psi d\tau \\ &= l(l+1) \hbar^2 \end{aligned} \quad (20)$$

Therefore

$$d = \frac{\lambda_c}{2\pi} \frac{l(l+1)}{n^2} \quad (21)$$

and

$$\epsilon = \left(1 - \frac{l(l+1)}{n^2} \right)^{1/2} \quad (22)$$

where the Compton wavelength is

$$\lambda_c = \frac{h}{mc} \quad (23)$$

It follows that:

$$\frac{L^2}{2} = \frac{2\pi}{\lambda_c} \alpha_{gs} \hbar^2 - (24)$$

and that
$$a = \frac{\lambda_c}{2\pi} \frac{n^2}{\alpha_{gs}} - (25)$$

Therefore Eq. (1) becomes:

$$\nabla^2 \psi = - \frac{2\pi}{\lambda_c} \alpha_{gs} \left(\frac{2}{r} - \frac{2\pi \alpha_{gs}}{\lambda_c n^2} \right) \psi - (26)$$

Finally note that:

$$\frac{2\pi}{\lambda_c} \alpha_{gs} = \frac{1}{a_0} = \frac{me^2}{4\pi \hbar^2 \epsilon_0} - (27)$$

where a_0 is the Bohr radius. So:

$$\boxed{\nabla^2 \psi = - \frac{1}{a_0} \left(\frac{2}{r} - \frac{1}{n^2 a_0} \right) \psi} - (28)$$

As shown in Note 371(8):

$$\nabla^2 \psi = \frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{d\psi(r)}{dr} \right) - \frac{l(l+1)}{r^2} \psi(r) - (29)$$

So the radial wavefunctions $\psi(r)$ can be obtained directly from eqs. (28) and (29). The angular wavefunctions are the spherical harmonics. The next note will extend this method to other atoms, notably helium.