

403(8): Final Version of Note 403(7)

Carries the solution of the ECE-2 force equation:

$$\frac{d^2}{dr^2} \left(\frac{1}{r} \right) + \frac{1}{r} = \frac{1}{d} (1 + r\omega_r) \quad - (1)$$

given by Maxima computer algebra:

$$\phi = - \int \left(\frac{d}{2\omega_r \log u - du^2 + 2u - 2C_1} \right)^{1/2} du \quad - (2)$$

where C_1 is a constant to be determined. Here:

$$u = \frac{1}{r}, \quad d = \frac{L^2}{m^2 m G} \quad - (3)$$

In the limit: $\omega_r \rightarrow 0 \quad - (4)$

Eq. (2) reduces to:

$$\phi = - \int \left(\frac{d}{-du^2 + 2u - 2C_1} \right)^{1/2} du \quad - (5)$$

which is the equation of a conic section, notably an ellipse, provided that:

$$\phi(r) = \int \frac{L}{r^2} \left(2m \left(H + \frac{mMG}{r} - \frac{L^2}{2mr^2} \right) \right)^{1/2} dr \quad - (6)$$

is the notation of Note 403(7).

Comparing eqs. (5) and (6):

$$\frac{L^2}{2m \left(H + \frac{mMG}{r} - \frac{L^2}{2mr^2} \right)} = \frac{d}{(-du^2 + 2u - 2C_1)} \quad - (7)$$

$$= \frac{L^2}{m^2 m_G} \left(\frac{1}{\frac{2mH}{m^2 m_G} + 2u - \frac{L^2 u^2}{m^2 m_G}} \right)$$

$$= \frac{d}{\frac{2H}{m m_G} + 2u - d u^2}$$

So the equations are the same if:

$$C_1 = -\frac{H}{m m_G} \quad - (8)$$

P.E.D.

The result (8) is also given by computer algebra.

The solution of eq. (6) is the conic section:

$$\frac{d}{r} = 1 + e \cos \phi \quad - (9)$$

and the half right distance is:

$$d = \frac{L^2}{m^2 m_G} \quad - (10)$$

and the eccentricity is:

$$e = \left(1 + \frac{2HL^2}{m^3 m^2 G^2} \right)^{1/2} \quad - (11)$$

The semi major axis is

$$a = \frac{d}{1 - e^2} = \frac{m m_G}{2|H|} \quad - (12)$$

so for eqs (7) and (12):

$$\frac{1}{a} = \frac{2|H|}{nm\bar{b}} \quad - (13)$$

So
$$\phi = - \int \left(\frac{d}{-du^2 + 2u + \frac{1}{a}} \right)^{1/2} du \quad - (14)$$

is the equation of a static ellipse, a is well known. This is changed by the spin correction to give:

$$\phi = - \int \left(\frac{d}{2\omega_r \log_e u - du^2 + 2u + \frac{1}{a}} \right)^{1/2} du \quad - (15)$$

Since ω_r is very small for small precessions, ω_r is a small perturbation of an ellipse. From note 405(6) it is known that this perturbation is the perihelion precession:

$$\Delta\phi = \frac{4}{3} \frac{\langle \underline{\delta r} \cdot \underline{\delta r} \rangle}{r^2} \quad - (16)$$

Now note that:

$$\begin{aligned} & 2\omega_r \log_e u - du^2 + 2u + \frac{1}{a} \\ &= \left(-du^2 + 2u + \frac{1}{a} \right) (1+x) \end{aligned} \quad - (17)$$

Here:

$$x = \frac{2\omega_r \log_e u}{-du^2 + 2u + \frac{1}{a}} \quad - (18)$$

So:

$$\phi = - \int \left(\frac{d}{\left(-du^2 + 2u + \frac{1}{a} \right) (1+x)} \right)^{1/2} du \quad - (19)$$

For small x :

$$(1+x)^{-1/2} \sim 1 - \frac{x}{2} \quad - (20)$$

$$\phi \sim - \int \left(\frac{d}{-du^2 + 2u + \frac{1}{a}} \right)^{1/2} \left(1 - \frac{x}{2} \right) du \quad - (21)$$

$$= - \int \left(\frac{d}{-du^2 + 2u + \frac{1}{a}} \right)^{1/2} + \frac{1}{2} \int x \left(\frac{d}{-du^2 + 2u + \frac{1}{a}} \right)^{1/2} du$$

Therefore the conversion to the ellipse is:

$$\Delta \phi = \frac{1}{2} \int x \left(\frac{d}{-du^2 + 2u + \frac{1}{a}} \right)^{1/2} du$$

$$= \int \frac{\omega_r d^{1/2} \log_e u}{\left(-du^2 + 2u + \frac{1}{a} \right)^{3/2}} du \quad - (22)$$

Now use:

$$\omega_r = \frac{2}{3} \frac{\langle \underline{sr} \cdot \underline{sr} \rangle}{r^3} = \frac{2}{3} \langle \underline{sr} \cdot \underline{sr} \rangle u^3 \quad - (23)$$

first let:

$$\Delta \phi = \frac{2}{3} \int \langle \underline{sr} \cdot \underline{sr} \rangle \left(\frac{u^3}{-du^2 + 2u + \frac{1}{a}} \right)^{3/2} \log_e u du \quad - (24)$$

If it is assumed that $\langle \underline{sr} \cdot \underline{sr} \rangle$ is approximately constant, then:

$$\Delta\phi \sim \frac{2}{3} \langle \underline{S}_r \cdot \underline{S}_r \rangle \int \left(\frac{u^2}{-d.u^2 + 2u + \frac{1}{a}} \right)^{3/2} \log_e u \, du$$

- (25)

at the perihelion:

$$u_{\min} = \frac{1+\epsilon}{d} = \frac{1}{a(1-\epsilon)} \quad - (26)$$

so from eq. (16)

$$\begin{aligned} \Delta\phi(\text{perihelion}) &= \frac{4}{3} \frac{\langle \underline{S}_r \cdot \underline{S}_r \rangle}{r_{\min}^2} \\ &= \frac{4}{3} \langle \underline{S}_r \cdot \underline{S}_r \rangle u_{\min}^2 \\ &= \frac{4}{3} \left(\frac{1+\epsilon}{d} \right)^2 \langle \underline{S}_r \cdot \underline{S}_r \rangle \quad - (27) \end{aligned}$$

Eq. (27) is a special case of eq. (25).

Maxima could be used to look for an analytical solution of eq. (25) or to integrate numerically. It is seen that the precession at perihelion is due to $\langle \underline{S}_r \cdot \underline{S}_r \rangle$.

Eq (25) has no solution in Wolfram but could be solved numerically or is an approximation.