

# 411(5): Comparison of the Einstein and Classical Theories of Precession.

As in previous UFT papers the Einstein theory is based on

$$\frac{d^2 u}{d\phi^2} + u = \frac{1}{d} + \delta u^2 \quad - (1)$$

where

$$u = \frac{1}{r}, \quad \frac{1}{d} = \frac{G m^2 M}{L^2}, \quad \delta = \frac{3 M G}{c^2} \quad - (2)$$

Computation in previous UFT paper has shown that eq. (1) fails catastrophically under well defined conditions, i.e. that  $m$  collides with  $M$ .

Eq. (1) can be solved analytically with a method given in Mars and Thornton, 3rd. edition, chapter 7.9. It is a method based on trial functions. The function:

$$u = u_1 = \frac{1}{d} (1 + \epsilon \cos \phi) \quad - (3)$$

produces the result:

$$\frac{d^2 u}{d\phi^2} + u = \frac{1}{d} + \frac{\delta}{d^2} (1 + \epsilon \cos \phi)^2 \quad - (4)$$

The function:

$$u_1 = \frac{1}{d} (1 + \epsilon \cos \phi) \quad - (5)$$

is a solution of:

$$\frac{d^2 u_1}{d\phi^2} + u_1 = \frac{1}{d} \quad - (6)$$

but is not a solution of eq. (1). To obtain a solution of eq. (1) use:

$$u = u_1 + u_p \quad - (7)$$

where  $u_p$  is a particular integral.

It is seen that:

$$\begin{aligned}\frac{d^2 u_1}{d\phi^2} + u_1 &= \frac{1}{d} + \frac{\delta}{d^2} (1 + 2\epsilon \cos \phi + \epsilon^2 \cos^2 \phi) \\ &= \frac{1}{d} + \frac{\delta}{d^2} \left( 1 + 2\epsilon \cos \phi + \frac{\epsilon^2}{2} (1 + \cos 2\phi) \right) \quad - (8)\end{aligned}$$

If  $u_p$  is chosen by inspection:

$$u_p = \frac{\delta}{d^2} \left[ 1 + \frac{\epsilon^2}{2} + \epsilon \phi \sin \phi - \frac{\epsilon^2}{6} \cos 2\phi \right] \quad - (9)$$

it follows that:

$$\frac{du_p}{d\phi} = \frac{\delta}{d^2} \left[ \epsilon \sin \phi + \epsilon \phi \cos \phi + \frac{\epsilon^2}{3} \sin 2\phi \right] \quad - (10)$$

and

$$\frac{d^2 u_p}{d\phi^2} = \frac{\delta}{d^2} \left[ \epsilon \cos \phi - \epsilon \phi \sin \phi + \epsilon \cos \phi + \frac{2}{3} \epsilon^2 \cos 2\phi \right] \quad - (11)$$

so

$$\begin{aligned}\frac{d^2 u_p}{d\phi^2} + u_p &= \frac{\delta}{d^2} \left( 1 + 2\epsilon \cos \phi + \frac{\epsilon^2}{2} (1 + \cos 2\phi) \right) \\ &= \frac{\delta}{d^2} (1 + \epsilon \cos \phi)^2 \quad - (12)\end{aligned}$$

So if  
then

$$u = u_1 + u_p \quad - (13)$$

$$\begin{aligned}\frac{d^2 u}{d\phi^2} + u &= \frac{1}{d} + \frac{\delta}{d^2} (1 + \epsilon \cos \phi)^2 \quad - (14) \\ &= \frac{1}{d} + \frac{\delta}{d^2} u_1^2\end{aligned}$$

3) This is an approximation to the result (1). The accurate form of  $z_v(14)$  is:

$$\frac{d^2}{d\phi^2}(u_1 + u_p) + u_1 + u_p = \frac{1}{d} + \delta(u_1 + u_p)^2 \quad (15)$$

This is approximated by:

$$\frac{d^2}{d\phi^2}(u_1 + u_p) + u_1 + u_p = \frac{1}{d} + \delta u_1^2 \quad (16)$$

So it is assumed that:

$$u_p \ll u_1 \quad (17)$$

This is true only for very small

$$\delta = \frac{3mG}{c^2} \quad (18)$$

and

$$\frac{\delta}{d^2} \ll 1 \quad (19)$$

This approximation:

$$u \sim u_1 + u_p = \left[ \frac{1}{d} (1 + \epsilon \cos \phi) + \frac{\delta \epsilon}{d^2} \phi \sin \phi \right] + \left[ \frac{\delta}{d^2} \left( 1 + \frac{\epsilon^2}{2} \right) - \frac{\delta \epsilon^2}{6d^2} \cos 2\phi \right] \quad (20)$$

This can be graphed by computer algebra and does not in general give a precise result.

4) The Euler theory collapses at this point, but special pleading is used as follows to try to rescue the situation. This dubious procedure is given on page 269 of the 1st edition and is reproduced as follows for the sake of argument. The first term in the second set of brackets of eq. (20) is a constant, and the second term "the bracket is regarded by Maria and Thornton as a 'small and periodic disturbance', so neither term gives precession. These claims can be checked by computer algebra for various  $\delta$  and  $d$ ."

The first set of brackets is described as the secular term:

$$u_{\text{secular}} = \frac{1}{d} \left( 1 + \epsilon \cos \phi + \frac{\delta \epsilon}{d} \phi \sin \phi \right) \quad - (21)$$

$$= \frac{1}{d} \left( 1 + \epsilon \cos \phi \right) + \frac{\delta \epsilon}{d^2} \phi \sin \phi$$

$$= u_1 + \frac{\delta \epsilon}{d^2} \phi \sin \phi$$

in the approximation:  $\delta/d^2 \ll 1 \quad - (22)$

The constant term:

$$u_c = \frac{\delta \epsilon}{d^2} \phi \sin \phi \quad - (23)$$

can be graphed using computer algebra. It can be seen by inspection that as:

$$\phi \rightarrow \infty \quad - (24)$$

5) then:  $u_c \rightarrow \infty - (25)$

because:  $0 \leq \sin \phi \leq 1 - (26)$

Therefore the Euler term collapses completely at the point because there is no upper bound on  $\phi$ .  
The upper bound of  $\sin \phi$  is 1.

Very special pleading is used by Marion and Thornton by using:

$$1 + \epsilon \cos\left(\phi - \frac{\delta}{\alpha} \phi\right) = 1 + \epsilon \left( \cos \phi \cos \frac{\delta}{\alpha} \phi + \sin \phi \sin \frac{\delta}{\alpha} \phi \right) \\ \sim 1 + \epsilon \cos \phi + \frac{\delta \epsilon}{\alpha} \phi \sin \phi - (27).$$

using the assumption that  $\delta$  is small, so:

$$\cos\left(\frac{\delta}{\alpha} \phi\right) \sim 1 - (28)$$

$$\sin\left(\frac{\delta}{\alpha} \phi\right) \sim \frac{\delta}{\alpha} \phi - (29)$$

Therefore the secular term becomes:

$$u_{\text{secular}} \sim \frac{1}{\alpha} \left( 1 + \epsilon \cos(\phi x) \right) - (30)$$

where  $x = 1 - \frac{\delta}{\alpha} - (31)$

The  $x$  substitution of ECE has developed eq. (30) in great detail. It gives a precessing orbit if and only if

$$b) \quad x - 1 \rightarrow 0 \quad - (32)$$

$$\text{i.e.} \quad s/d \ll 1 \quad - (33)$$

otherwise eq. (31) gives very irregular orbits  
which are never observed experimentally  
 At this point the Einstein orbital theory collapses

again. However Maria and Thornton ignore all these  
 flaws, and proceed by arguing that for one orbit:

$$\phi - \frac{s}{d} \phi = 2\pi \quad - (34)$$

So

$$\phi = \frac{2\pi}{1 - s/d} \sim 2\pi \left(1 + \frac{s}{d}\right) \quad - (35)$$

under condition (33).

The precession is therefore:

$$\Delta \phi = 2\pi \frac{s}{d} \quad - (36)$$

where

$$x = \frac{s}{d} = \frac{3MG}{c^2 d} \quad - (37)$$

Finally Maria and Thornton claim without any  
 justification (see 4FT410) that this theory is always  
 very precise.

No rational or enlightened thinker would  
accept such a theory. The most devastating  
 criticism of it is that

7)  $u_{\text{scaler}} = u_1 + \frac{\delta E}{d^2} \phi \sin \phi - (38)$

$\rightarrow \infty$

so  $\phi \rightarrow \infty - (39)$

and after sufficient revolutions:  
 $\frac{\delta E}{d^2} \phi \sin \phi \gg u_1 - (39.40)$

and the orbit ceases completely to be a circle.  
section. This is a rotation of disc.

This disaster for EBR can be easily illustrated by complex algebra.

The Classical Theory of Precession

As shown in the previous note, the frame rotation:

$\phi' = \phi + \omega t - (41)$

produces the orbit:

$r = \frac{d'}{1 + \epsilon' \cos(\phi + \omega_1 t)} - (42)$

At the perihelion:

$\phi + \omega_1 t = 2\pi - (43)$

so

$\phi = 2\pi - \omega_1 t - (44)$

The precession is defined by:

$\Delta \phi = 2\pi - (2\pi - \omega_1 t) - (45)$   
 $= \omega_1 t$

8) If the time taken for one orbit is  $T$ , then after one orbit:

$$\Delta\phi = \omega_1 T \quad - (46)$$

So any precession is described by  $\omega_1$ , the angular velocity of frame rotation.

This is a far simpler and much more elegant theory than the Einstein theory, and contains no errors. The orbit shrinks by:

$$r^2 = \frac{L'}{m(\omega + \omega_1 + T \frac{d\omega_1}{dt})} \quad - (47)$$

where  $\omega$  is the orbital angular velocity and  $\omega_1$  is the angular velocity of frame rotation. The time  $T$  is found from Kepler's 3rd law.

The ECE2 covariant unified field theory is needed to show that the origin of  $\omega_1$  is spacetime torsion. This is missing completely from the Einstein theory.

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